

# First Measurement of the Fraction of Top Quark Pair Production Through Gluon-Gluon Fusion

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We present the first measurement of  $\sigma(gg \rightarrow t\bar{t})/\sigma(p\bar{p} \rightarrow t\bar{t})$ . We use  $0.96 \text{ fb}^{-1}$  of  $\sqrt{s} = 1.96 \text{ TeV}$   $p\bar{p}$  collision data recorded with the CDF II detector at Fermilab. We identify the candidate  $t\bar{t}$  events with a high-energy charged lepton, a neutrino candidate, and four or more jets. Using charged particles with low transverse momentum in  $t\bar{t}$  events, we find  $\sigma(gg \rightarrow t\bar{t})/\sigma(p\bar{p} \rightarrow t\bar{t}) = 0.07 \pm 0.14(\text{stat}) \pm 0.07(\text{syst})$ , corresponding to a 95% confidence level upper limit of 0.33, in agreement with the standard model NLO prediction of  $0.15 \pm 0.05$ .

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Many studies have been dedicated to the understanding of the top quark, motivated in part by its large mass that may give it a unique role in the generation of mass for the quarks, leptons, and force carriers in the standard model (SM) of particle physics. In  $p\bar{p}$  collisions at a center-of-momentum en-

ergy of  $\sqrt{s} = 1.96 \text{ TeV}$ ,  $(15 \pm 5)\%$  of  $t\bar{t}$  pairs are expected to be produced through gluon-gluon fusion and the rest through quark-antiquark annihilation [1, 2], based on next-to-leading-order (NLO) quantum chromodynamics (QCD) calculations. The inclusive  $t\bar{t}$  production cross section has been measured by both CDF [3, 4] and D0 [5] collaborations using various methods and decay modes of the  $t\bar{t}$  pairs, and the results are in agreement with SM predictions. However, the details of the production process have never been investigated.

A measurement of the  $gg \rightarrow t\bar{t}$  production cross section tests the QCD prediction. Also of interest is any indication of new top quark production and decay mechanisms. As the partonic cross section calculations are directly related to the momentum distributions of constituents of the colliding protons [1], such a measurement could assist in reducing the uncertainties in the gluon distributions within protons.

In this Letter we report the first measurement of the fractional cross section of  $t\bar{t}$  production through gluon-gluon fusion,  $\sigma(gg \rightarrow t\bar{t})/\sigma(p\bar{p} \rightarrow t\bar{t})$ . To discriminate between the

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similar final state signatures of  $gg \rightarrow t\bar{t}$  and  $q\bar{q} \rightarrow t\bar{t}$ , we take advantage of the higher probability for a gluon than for a quark to radiate a low-momentum gluon [6]. Therefore, on average we expect a larger low-momentum charged particle multiplicity in  $gg \rightarrow t\bar{t}$  compared to  $q\bar{q} \rightarrow t\bar{t}$ . Given the large theoretical uncertainties associated with gluon radiation, we do not rely on theoretical calculations for the modeling of the charged particle multiplicity. Instead, we use two different processes,  $W+n$  jet and two-jet (dijet) production, with well-understood production mechanisms, as calibration samples to relate the observed charged particle multiplicity to the fraction of processes involving more gluons.

We use a data sample of  $\sqrt{s} = 1.96$  TeV  $p\bar{p}$  collisions with an integrated luminosity of  $0.96 \pm 0.06 \text{ fb}^{-1}$  recorded by the CDF II detector at Fermilab between March 2002 and February 2006. The CDF II detector is described in detail in [7]; here, we briefly discuss the components essential for this analysis. The detector consists of a tracking system immersed in a solenoidal magnetic field of 1.4 T and electromagnetic and hadronic calorimeters surrounding the solenoid, followed by the muon system. Electrons, photons, and hadronic jets are identified using calorimeters and the tracking information. Muons are identified by the muon system together with tracking and calorimeter information. The data are collected using a three-level trigger system.

According to the SM top quarks almost always decay to a  $W$  boson and a bottom quark, and so in  $t\bar{t}$  events we expect to have two  $W$  bosons and two  $b$  quarks. We select  $t\bar{t}$  candidate events where one of the  $W$  bosons decays to two jets and the other decays to a lepton ( $l$ ) and the corresponding neutrino. In this analysis  $l$  is either an electron or a muon. Our first calibration dataset is a set of  $W(\rightarrow l\nu)+n$  jet ( $n = 0, 1, 2, 3$ ) candidate events, for which the number of gluons involved in the production process increases with the number of jets [8]. The second is a set of events with two back-to-back, high-energy jets. The average number of gluons involved in dijet production [9] falls with increasing transverse energy ( $E_T$ ) [10] of the highest  $E_T$  jet (leading jet), as the relative rate of the  $q\bar{q} \rightarrow q\bar{q}$ ,  $q\bar{q} \rightarrow qg$  and  $gg \rightarrow gg$  subprocesses change. The number of gluons in each subprocess is 0, 2, and 4 respectively.

The  $W+n$  jet data are collected with an inclusive lepton trigger that requires an electron with  $E_T > 18$  GeV or a muon with  $p_T > 18$  GeV/ $c$ . We select events with a reconstructed isolated electron (muon) candidate with  $E_T > 20$  GeV ( $p_T > 20$  GeV/ $c$ ) and a missing  $E_T$  ( $\cancel{E}_T$ )  $> 20$  GeV. We categorize the  $W+n$  jet samples by  $n$ , the number of jet candidates with  $E_T > 15$  GeV and pseudorapidity region  $|\eta| < 2$ . Jets are defined using an iterative cone algorithm [11] with a cone of  $\Delta R = \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2} = 0.4$  and are corrected for absolute energy response,  $\eta$  dependence of calorimeter response, and multiple interactions. For the  $t\bar{t}$  data sample, in addition to the above, we require four or more jets where at least one is identified as originating from a  $b$  quark ( $b$  tag). To define a  $b$  tag, we identify within a jet a long-lived  $B$ -hadron candidate through the presence of a displaced secondary vertex [3]. In both  $t\bar{t}$  and  $W+n$  jet samples, we remove any event

with a second lepton candidate consistent with arising from a  $Z$  boson decay or a  $t\bar{t}$  event in which both  $W$  bosons decay to leptons. We also veto the events in which the electron (muon) is consistent with coming from a conversion photon (cosmic ray) [3]. The dijet data are collected using two inclusive jet triggers that require a jet with  $E_T$  of at least 50 GeV or at least 100 GeV (Jet50 and Jet100 datasets). We require a minimum leading jet  $E_T$  of 75 and 130 GeV for Jet50 and Jet100 datasets, respectively, to avoid any trigger bias. We remove events containing an electron (muon) candidate with  $E_T > 20$  GeV ( $p_T > 20$  GeV/ $c$ ). We also require exactly two jets with  $|\eta| \leq 2$  and a minimum  $E_T$  of 20 GeV and with the two jets back-to-back, having a  $|\Delta\phi| \geq 2.53$  rad.

The background processes in our  $t\bar{t}$  sample consist of  $W+$  jets, electroweak processes ( $WW$ ,  $WZ$ ,  $ZZ$ ), single top quark, and multi-jet QCD processes (non- $W$ ). For non- $W$  and  $W+$  jets background, we can have a real  $b$  tag (heavy flavor background, HF) or have a  $b$  tag due to misidentification (light flavor background, LF). We estimate LF and HF in events with a real  $W$  boson using various calibration datasets. For the small fraction of events from non- $W$  sources, we assume the non- $W$  background is equal parts HF and LF. The results of the analysis are insensitive to this assumption. Single top quark processes are part of HF, while diboson backgrounds, ignoring the few  $Z \rightarrow b\bar{b}$  events, are included in LF. We find 240  $t\bar{t}$  candidates with an estimated background contamination of  $(13 \pm 2)\%$ . The background estimates are found using the method explained in [3].

The number of low- $p_T$  charged particles  $N_{\text{trk}}$  is affected by low-energy particles arising from jet fragmentation as well as multiple interactions within the same  $p\bar{p}$  bunch crossing. To include a track in our definition of  $N_{\text{trk}}$ , we require it to have a  $p_T$  in the range 0.3-2.9 GeV/ $c$  and  $|\eta| \leq 1.1$  and to originate from the vertex associated with the charged leptons and jets. We reject the track if it falls within  $\Delta R = 0.6$  and  $\Delta R = 0.4$  of jets with  $E_T \geq 15$  GeV (high  $E_T$  jets) and  $6 \leq E_T < 15$  GeV (low  $E_T$  jets), respectively. Excluding these tracks results in different available tracking area for each event. We therefore correct the observed multiplicity to the total tracking coverage in  $\eta$  and  $\phi$  event-by-event. The resulting track multiplicity still has a modest dependence on the number of high  $E_T$  jets in the event. We therefore make a further correction to  $N_{\text{trk}}$  by measuring this dependence in multi-jet QCD candidate events and using this as a per-jet correction ( $\sim 1$  track per jet) to the multiplicity for all jets with  $|\eta| \leq 1.1$ .

We show that there is a correlation between the average number of low- $p_T$  charged particles  $\langle N_{\text{trk}} \rangle$  and the average number of gluons involved in the production process  $\langle N_g \rangle$  in a given sample. We count the number of gluons that are part of the production process using Monte Carlo (MC) calculations for both dijet and the  $W+n$  jet data samples. The  $W+n$  jet MC sample is created using the ALPGEN [12] program followed by PYTHIA [13] to perform the jet fragmentation. The MC dijet events are created using the PYTHIA MC. We plot the observed  $\langle N_{\text{trk}} \rangle$  in data against the expected  $\langle N_g \rangle$  from MC calculations for the calibration samples in Fig. 1. This

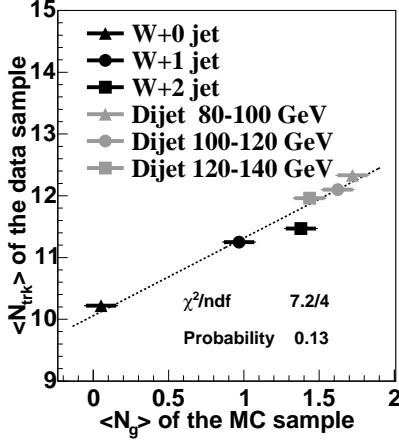


FIG. 1: The correlation between the average number of low- $p_T$  charged particles (data) and the average number of gluons (MC). The dotted line is from a linear fit to the points.

demonstrates an approximately linear dependence between  $\langle N_{\text{trk}} \rangle$  and  $\langle N_g \rangle$ . We do not use this plot to obtain our result, but rather directly fit the observed  $N_{\text{trk}}$  distributions as described below.

The  $\langle N_{\text{trk}} \rangle$  and  $\langle N_g \rangle$  correlation enables us to define  $N_{\text{trk}}$  distributions each representing a specific average number of gluons involved in the production process. We use this correlation and the observed  $N_{\text{trk}}$  distributions in the  $W + 0$  jet sample and the dijet sample with leading jet  $E_T$  of 80-100 GeV to define a no-gluon and a gluon-rich  $N_{\text{trk}}$  distribution, respectively. To do so, an iterative procedure is adopted in order to remove the no-gluon (gluon-rich) contribution from the  $N_{\text{trk}}$  distribution of the 80-100 GeV dijet ( $W + 0$  jet) sample. There are no significant changes in the distributions after the first iteration. The  $W + 0$  jet sample is largely composed of the Drell-Yan  $q\bar{q}'$  process with a small QCD background of order 4% and contribution from  $W$  production in association with other partons where none of the final state jets are detected. The fraction of  $W + 0$  jet candidates with production processes involving gluons is estimated to be  $(5 \pm 4)\%$ . The no-gluon contribution of dijet candidates with leading jet  $E_T$  of 80-100 GeV comes from  $qq \rightarrow qq$  processes and is estimated to be  $(27 \pm 3)\%$ .

To verify that the no-gluon or gluon-rich distribution can model the  $N_{\text{trk}}$  distribution of any process with comparable  $\langle N_g \rangle$  regardless of the center-of-momentum energy, we check the  $W + 1$  jet data sample, and we see no dependence on jet  $E_T$  in  $\langle N_{\text{trk}} \rangle$ .

The gluon-rich fraction associated with a given  $N_{\text{trk}}$  distribution can be found using a binned likelihood fit of the  $N_{\text{trk}}$  distribution of the form

$$N[f_g \mathcal{F}_g(N_{\text{trk}}) + (1 - f_g) \mathcal{F}_q(N_{\text{trk}})], \quad (1)$$

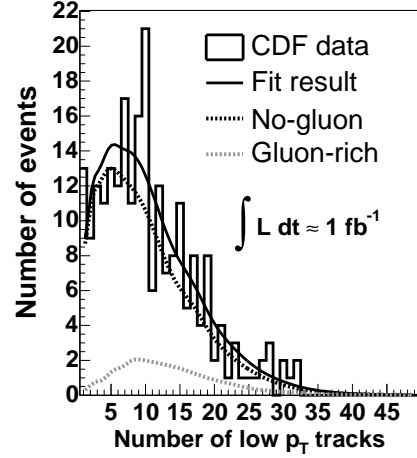


FIG. 2: The number of low- $p_T$  charged particles for the  $t\bar{t}$  candidates, the fit result, the gluon-rich, and no-gluon components.

where  $N$  is the normalization factor and one of the free parameters,  $f_g$  is the fraction of gluon-rich components of the sample and the second free parameter, and  $\mathcal{F}_g(N_{\text{trk}})$  and  $\mathcal{F}_q(N_{\text{trk}})$  are the normalized gluon-rich and no-gluon parameterizations, respectively. We have verified this technique is free of bias using randomly generated samples from the data  $N_{\text{trk}}$  distributions.

The  $N_{\text{trk}}$  distribution of the  $t\bar{t}$  candidates, shown in Fig. 2, has a mean of  $10.6 \pm 0.5$ . The fit, shown in the figure, models the data distribution very well, based on a goodness of fit test with 92% probability. The measured gluon-rich fraction in  $t\bar{t}$  candidates determined by fitting the  $N_{\text{trk}}$  distribution consists of two components, the  $t\bar{t}$  gluon-rich fraction and the background gluon-rich fraction. Therefore, knowing the background fraction in our sample  $f_b$  and the measured  $f_g$  from the fit, we can write

$$f_g = f_b f_g^{\text{bkg}} + (1 - f_b) f_g^{t\bar{t}}, \quad (2)$$

where  $f_g^{\text{bkg}}$  and  $f_g^{t\bar{t}}$  are the gluon-rich fraction of the background and  $t\bar{t}$  signal, respectively. To estimate  $f_g^{\text{bkg}}$ , we measure  $f_g$  in the  $W + 1$ ,  $W + 2$ , and  $W + 3$  jet data samples with no  $b$  tag and with at least one  $b$  tag using the fit to the  $N_{\text{trk}}$  distribution for each of these samples. We then extrapolate the  $f_g$  values from the  $W + 1$ , 2, and 3 jet samples to  $W + 4$  or more jet bins for both  $b$ -tag and no  $b$ -tag samples. We consider the  $b$ -tag sample as representative of HF and the no  $b$ -tag sample as representative of LF. Using these extrapolations and the LF and HF fractions, we find  $f_g^{\text{bkg}} = 0.54 \pm 0.09$  and  $f_g^{t\bar{t}} = 0.09 \pm 0.16$ .

Given  $f_g^{t\bar{t}}$ , we measure  $\sigma(gg \rightarrow t\bar{t})/\sigma(p\bar{p} \rightarrow t\bar{t})$  as

$$\left[ 1 - \frac{\mathcal{A}_{gg}}{\mathcal{A}_{q\bar{q}}} + \left( \frac{\mathcal{A}_{gg}}{\mathcal{A}_{q\bar{q}}} \right) \left( \frac{1}{f_{gtt}^{tt}} \right) \right]^{-1} = 0.07 \pm 0.14(\text{stat}), \quad (3)$$

where  $\mathcal{A}_{gg}$  and  $\mathcal{A}_{q\bar{q}}$  are the acceptance for  $gg \rightarrow t\bar{t}$  and  $q\bar{q} \rightarrow t\bar{t}$ , respectively. Using PYTHIA [13] MC calculations, we find  $(14.1 \pm 0.5)\%$  and  $(11.5 \pm 0.4)\%$  for  $\mathcal{A}_{gg}$  and  $\mathcal{A}_{q\bar{q}}$ , respectively. The acceptance uncertainties include the systematic uncertainties.

The systematic uncertainties of this measurement, a total of 0.07, arise from uncertainties in the measurement of  $N_{\text{trk}}$  and the subsequent calculations. The uncertainties in  $N_{\text{trk}}$  are due to the per-jet correction(0.05), the estimated gluon content of the  $W+0$  jet sample(0.04), and the choice of the low  $E_T$  jet cut(0.02). In addition to these sources, there are uncertainties associated with the estimated  $qq \rightarrow qq$  fraction of the 80-100 GeV dijet sample, the background fraction, the modeling of the background gluon-rich fraction, the non- $W$  background fraction, and the acceptances; these are all negligible. To estimate the effects of all the above uncertainties, we changed the central values and measured the change in the relevant variables.

The result corresponds to an upper limit of 0.33 at 95% confidence level. We use a classical statistical technique to set the limit by simulating the possible outcomes for a given true value taking into account the systematic effects.

In conclusion, we have presented the first measurement of  $\sigma(gg \rightarrow t\bar{t})/\sigma(p\bar{p} \rightarrow t\bar{t})$  and found  $0.07 \pm 0.14(\text{stat}) \pm 0.07(\text{syst})$ , corresponding to an upper limit of 0.33 at 95% confidence level, in  $0.96 \text{ fb}^{-1}$  of data collected at CDF. This is in agreement with the SM prediction of  $0.15 \pm 0.05$ , and does not suggest that non-standard model processes [14] contribute to top quark pair production at the Tevatron.

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Korean Science and Engineering Foundation and the Korean Research Foundation; the Science and Technology Facilities Council and the Royal Society, UK; the Institut National de Physique Nucleaire et Physique des Particules/CNRS; the Russian Foundation for Basic Research; the Comisión Interministerial de Ciencia y Tecnología, Spain; the European Community's Human Potential Programme; the Slovak R&D Agency; and the Academy of Finland.

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